

## Lecture 14

Last time:  $w(t, \tau) \Rightarrow w(t - \tau)$

$$y(t) = \int_{-\infty}^t w(t - \tau)x(\tau)d\tau$$

Let:  $\begin{cases} \tau' = t - \tau \\ -d\tau = d\tau' \end{cases}$

$$y(t) = \int_0^\infty w(\tau')x(t - \tau')d\tau'$$

For the differential system characterized by its equations of state, specialization to invariance means that the system matrices  $A, B, C$  are constants.

$$\begin{aligned} \dot{\underline{x}} &= A\underline{x} + B\underline{u} \\ \underline{y} &= C\underline{x} \end{aligned}$$

For  $A, B, C$  constant:

$$\underline{y}(t) = C\underline{x}(t)$$

$$\underline{x}(t) = \Phi(t - t_0)\underline{x}(t_0) + \int_{t_0}^t \Phi(t - \tau)B\underline{u}(\tau)d\tau$$

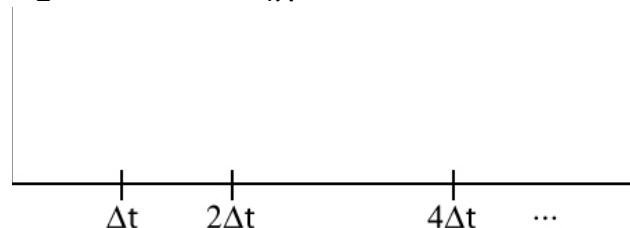
The transition matrix can be expressed analytically in this case.

$$\frac{d}{dt}\Phi(t, \tau) = A\Phi(t, \tau), \quad \text{where } \Phi(\tau, \tau) = I$$

This is a matrix form of first order, constant coefficient differential equation. The solution is the matrix exponential.

$$\Phi(t, \tau) = e^{A(t-\tau)}$$

$$e^{A(t-\tau)} = I + A(t - \tau) + \frac{1}{2}A^2(t - \tau)^2 + \dots + \frac{1}{k!}A^k(t - \tau)^k + \dots$$



Useful for computing  $\Phi(t)$  for small enough  $t - \tau$ .

The solution is

$$\underline{x}(t) = C\underline{x}(t)$$

$$\underline{x}(t) = e^{A(t-t_0)} \underline{x}(t_0) + \int_{t_0}^t e^{A(t-\tau)} B \underline{u}(\tau) d\tau$$

For  $t_0 \rightarrow \infty$ :

$$\begin{aligned} \underline{x}(t) &= \int_{-\infty}^t e^{A(t-\tau)} B \underline{u}(\tau) d\tau \\ &= \int_0^{\infty} e^{A\tau'} B \underline{u}(t-\tau') d\tau' \end{aligned}$$

and for a single input, single output (SISO) system,

$$w(t) = \underline{c}^T e^{At} \underline{b}$$

If  $x(t) = e^{j\omega t}$  for all past time

$$\begin{aligned} y(t) &= \int_0^{\infty} w(\tau) e^{j\omega(t-\tau)} d\tau \\ &= \left[ \int_0^{\infty} w(\tau) e^{-j\omega\tau} d\tau \right] e^{j\omega t} \\ &= F(\omega) x(t) \end{aligned}$$

Since  $w(\tau) = 0$  for  $\tau < 0$  for a realizable system, we see that the *steady state sinusoidal response function*,  $F(\omega)$ , for a system is the Fourier transform of the weighting function – where the coefficient unity must be used.

$$F(\omega) = \int_{-\infty}^{\infty} w(\tau) e^{-j\omega\tau} d\tau$$

and  $w(\tau)$  for a stable system is Fourier transformable.

Then

$$w(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} F(\omega) e^{j\omega t} d\omega$$

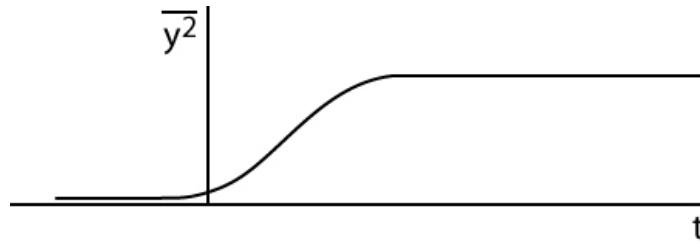
You can compute the response to any input at all, including transient responses, having defined  $F(\omega)$  for all frequencies.

The *static sensitivity* of the system is the zero frequency gain,  $F(0)$ , which is just the integral of the weighting function.

$$F(0) = \int_0^{\infty} w(\tau) d\tau$$

### **Stationary statistics**

Invariant output statistics implies more than stationary inputs and invariant systems; it also implies that the system has been in operation long enough under the present conditions to have exhausted all transients.



### **Input-Output Relations for Correlation and Spectral Density Functions**

Derive autocorrelation of output in terms of autocorrelation of input

$$\begin{aligned}
 y(t) &= \int_0^\infty w(\tau_1)x(t - \tau_1)d\tau_1 \\
 \bar{y} &= \int_0^\infty w(\tau_1)\bar{x}d\tau_1 \\
 &= \bar{x}\int_0^\infty w(\tau_1)d\tau_1 \\
 R_{yy}(\tau) &= \overline{y(t)y(t + \tau)} \\
 &= \int_0^\infty d\tau_1 w(\tau_1) \int_0^\infty d\tau_2 w(\tau_2) \overline{x(t - \tau_1)x(t + \tau - \tau_2)} \\
 &= \int_0^\infty d\tau_1 w(\tau_1) \int_0^\infty d\tau_2 w(\tau_2) R_{xx}(\tau + \tau_1 - \tau_2)
 \end{aligned}$$

$$\begin{aligned}
 \bar{y^2} &= \int_0^\infty d\tau_1 w(\tau_1) \int_0^\infty d\tau_2 w(\tau_2) R_{xx}(\tau_1 - \tau_2) \\
 R_{xy}(\tau) &= \overline{x(t)y(t + \tau)} \\
 &= \overline{x(t)\int_0^\infty w(\tau_1)x(t + \tau - \tau_1)d\tau_1} \\
 &= \int_0^\infty w(\tau_1)R_{xx}(\tau - \tau_1)d\tau_1
 \end{aligned}$$

Transform to get power density spectrum of output.

$$\begin{aligned}
 \bar{y} &= \bar{x} \int_0^\infty w(\tau) d\tau \\
 &= F(0)\bar{x} \\
 S_{yy}(\omega) &= \int_{-\infty}^{\infty} R_{yy}(\tau) e^{-j\omega\tau} d\tau \\
 &= \int_{-\infty}^{\infty} d\tau \int_0^\infty d\tau_1 w(\tau_1) \int_0^\infty d\tau_2 w(\tau_2) R_{xx}(\tau + \tau_1 - \tau_2) e^{-j\omega\tau_1} \\
 &= \underbrace{\int_{-\infty}^{\infty} d\tau R_{xx}(\tau + \tau_1 - \tau_2) e^{-j\omega(\tau+\tau_1-\tau_2)} \int_0^\infty d\tau_1 w(\tau_1) e^{j\omega\tau_1}}_{\text{first integral}} \int_0^\infty d\tau_2 w(\tau_2) e^{-j\omega\tau_2} \\
 \text{In first integral only, let } &\begin{cases} \tau' = \tau + \tau_1 - \tau_2 \\ d\tau' = d\tau \end{cases} \\
 S_{yy}(\omega) &= \int_{-\infty}^{\infty} d\tau' R_{xx}(\tau') e^{-j\omega\tau'} \int_{-\infty}^{\infty} d\tau_1 w(\tau_1) e^{j\omega\tau_1} \int_{-\infty}^{\infty} d\tau_2 w(\tau_2) e^{-j\omega\tau_2} \\
 &= S_{xx}(\omega)F(-\omega)F(\omega) \\
 &= |F(\omega)|^2 S_{xx}(\omega)
 \end{aligned}$$

The power spectral density thus does not depend upon phase properties.

The *cross-spectral density function* can be derived similarly, to obtain:

$$S_{xy}(\omega) = F(\omega)S_{xx}(\omega)$$

Mean squared output in time and frequency domain

$$\begin{aligned}
 \bar{y^2} &= R_{yy}(0) = \int_0^\infty d\tau_1 w(\tau_1) \int_0^\infty d\tau_2 w(\tau_2) R_{xx}(\tau_1 - \tau_2) \\
 &= \frac{1}{2\pi} \int_{-\infty}^{\infty} S_{yy}(\omega) d\omega \\
 &= \frac{1}{2\pi} \int_{-\infty}^{\infty} F(\omega)F(-\omega)S_{xx}(\omega) d\omega
 \end{aligned}$$

Generally speaking, with linear invariant systems it is easier to work in the transform domain than the time domain – so we shall commonly use the last expression to calculate the mean squared output of a system. However, control engineers are more accustomed to working with Laplace transforms than with Fourier transforms. By making the change of variables  $s = j\omega$  we can cast this expression in that form.

$$\begin{aligned}\overline{y^2} &= \frac{1}{2\pi} \int_{-\infty}^{j\infty} F\left(\frac{s}{j}\right) F\left(-\frac{s}{j}\right) S_{xx}\left(\frac{s}{j}\right) \frac{ds}{j} \\ &= \frac{1}{2j\pi} \int_{-\infty}^{\infty} F'(s) F'(-s) S'_{xx}(s) ds\end{aligned}$$

We know that  $S_{xx}(\omega)$  is even. If it is a rational function of  $\omega$ , and we will work exclusively with rational spectra, it is then a rational function of  $\omega^2$ . So only even powers of  $\omega$  appear in  $S_{xx}(\omega)$  and thus  $S_{xx}\left(\frac{s}{j}\right)$  which we may call  $S_{xx}(s)$  is derived from  $S_{xx}(\omega)$  by replacing  $\omega^2$  by  $-s^2$ .

$F'(s)$  is the ordinary transfer function of the system – the Laplace transform of its weighting function. Because  $w(t) = 0$ ,  $t < 0$ .

We shall drop the primes from now on.

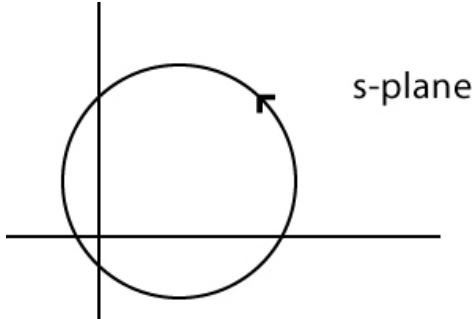
$$\overline{y^2} = \frac{1}{2\pi j} \int_{-\infty}^{j\infty} F(s) F(-s) S_{xx}(s) ds$$

$$\left. \begin{array}{l} \omega^2 = -s^2 \\ \omega^4 = s^4 \end{array} \right\} \text{in } S_{xx}(s)$$

### *Integrating the output spectrum*

#### General method

Cauchy Residue Theorem



$$\oint_C F(s) ds = 2\pi j \sum \text{(residues of } F(s) \text{ at the poles enclosed in the contour } C\text{)}$$

If  $F(s)$  has a pole of order  $m$  at  $s=a$ ,

$$\text{Res}(a) = \frac{1}{(m-1)!} \left\{ \frac{d^{m-1}}{ds^{m-1}} \left[ (s-a)^m F(s) \right]_{s=a} \right\}$$

$F(s)$  has a pole of order  $m$  at  $s=a$  if  $m$  is the smallest integer for which

$$\lim_{s \rightarrow a} \left[ (s-a)^m F(s) \right]$$

is finite.

If  $F(s)$  is rational and has a 1<sup>st</sup> order pole at  $a$ ,

$$F(s) = \frac{N(s)}{D(s)}$$
$$= \frac{N(s)}{(s-a)(s-b)\dots}$$

then

$$\text{Res}(a)_{m=1} = \lim_{s \rightarrow a} \left[ (s-a)F(s) \right]$$
$$= \frac{N(a)}{(a-b)(a-c)\dots}$$