## Strain and single-crystal elasticity

## **Assigned Reading:**

Nye JF (1957) Physical Properties of Crystals. Oxford University Press, Oxford, UK (Chapters 5 and 6).

## Resource reading (If you would like to brush up on stress, strain, and elasticity):

Means, W.D., *Stress and Strain*, 338pp pp., Springer-Verlag, New York, 1976. . (Many problems on which to practice.)

## **Infinitesimal Strain**



One dimensional distortion

Consider a one-dimensional object that is stretched. It is easy to imagine that the tail of the line segment,  $\Delta x$ , will move by an amount u, and that the tip of the line segment will be moved to  $u + \Delta x + \Delta u$ , where u is the rigid body translation and  $\Delta u$  is the amount that the length of the line segment has increased.



Strain at a point P is the increase in length/original length

 $\varepsilon = \frac{\Delta u}{\Delta x}$  and the total length of the stretched line segment  $\overline{P'Q'}$  is  $\Delta x + \Delta u = \Delta x (1 + \varepsilon)$ 

Goal: For a given arbitrary deformation of a body, define a tensor that describes the change of direction and length of any vector in the undeformed body.

Question: Why must the object be a tensor? What rank?



Point *P* is carried to *P*'; Point *Q* is carried to *Q*'. Since we are not interested in the rigid body translation, the deformation can be characterized by the coefficients  $\Delta u_i$ .

For any continuous differentiable variable,

$$\Delta u_i = \frac{\partial u_i}{\partial x_i} \Delta x_j$$

and we can define the displacement gradient tensor

$$e_{ij} = \frac{\partial u_i}{\partial x_j}$$

<u>Comments and Questions:</u> Because the tensor is written as a differential we are implicitely assuming small deformations (hence infinitesimal deformation gradient tensor).

e<sub>ii</sub> is a second rank tensor. How do you know?

As an exercise, explicitly write out the expression for the  $\Delta u_i$  using the deformation gradient tensor and  $\Delta x_i$ .

Now divide the tensor into two parts:

$$e_{ij} = \frac{1}{2} (e_{ij} + e_{ji}) + \frac{1}{2} (e_{ij} - e_{ji}) = \boxed{\varepsilon_{ij} + \omega_{ij} \equiv e_{ij}}$$

Where  $\varepsilon_{ij}$ ,  $\omega_{ij}$  are the infinitesimal strain and rotation tensors, respectively. The former is symmetric and the latter antisymmetric.