3.23 Electrical, Optical, and Magnetic Properties of Materials Fall 2007

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#### Last time: Wave mechanics

- 1. Time-dependent Schrödinger equation
- 2. Separation of variables stationary Schrödinger equation
- 3. Wavefunctions and what to expect from them
- 4. Free particle and particle in a 1-d, 2-d, 3-d box
- 5. Scanning tunnelling microscope
- 6. (Applets)

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#### Good news

 Study material: Prof Fink QM notes (uploaded on Stellar)

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#### First postulate

- All information of an ensemble of identical physical systems is contained in the ket |Ψ⟩ (usually a wavefunction Ψ(x,y,z,t), which is complex, continuous, finite, and single-valued, square-integrable (i.e. [||Ψ||<sup>2</sup> dr is finite)
- The ket can also be a geometrical vector (e.g. spin); in truth, wavefunctions are objects that satisfy vector algebra, and the space of wavefuncitons is a Hilbert space (instead of being 3-d, it's infinite-d)

#### Normalization, scalar products

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#### Second Postulate

 For every physical observable there is a corresponding Hermitian operator

#### From classical mechanics to operators

- Total energy is T+V (Hamiltonian is kinetic + potential)  $T = \frac{\rho^2}{2m}$
- classical momentum  $\vec{p} \rightarrow \rightarrow$  $\rightarrow$  gradient operator  $-i\hbar\vec{\nabla}$
- classical position  $\vec{r} \rightarrow$  $\rightarrow$  multiplicative operator  $\hat{r}$

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#### **Operators and operator algebra**

• Examples: derivative, multiplicative



#### Linear and Hermitian

• 
$$\hat{A}[\alpha|\varphi\rangle + \beta|\psi\rangle] - \alpha\hat{A}|\varphi\rangle + \beta\hat{A}|\psi\rangle$$
  
 $\vec{\nabla} \left[\alpha \varphi(\vec{r}) + \beta \varphi(\vec{r})\right] =$   
 $= \vec{\nabla} \left[\alpha \varphi(\vec{r})\right] + \vec{\nabla} \left[\beta \varphi(\vec{r})\right] =$   
 $= \alpha \nabla \varphi(\vec{r}) + \beta \nabla \psi(\vec{r})$   
 $= \alpha \nabla \varphi(\vec{r}) + \beta \nabla \psi(\vec{r})$   
 $\int \varphi^{\dagger}(\vec{r}) \left(\hat{A} \psi(\vec{r})\right) d\vec{r} = \int \left(\hat{A} \varphi(\vec{r})\right)^{\dagger} \varphi(\vec{r}) d\vec{r}$   
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Examples: 
$$(d/dx)$$
 and  $i(d/dx)$   
 $\zeta \varphi [A \varphi] = \zeta A \varphi [\varphi]$   
 $\int \varphi d \varphi d n = \int (d \varphi)^* \varphi d n$   
 $\int (d \varphi)^* \varphi d n = \int (d \varphi)^* \varphi d n$   
 $\int (\varphi \varphi) = (\varphi d \varphi) + f \varphi d \varphi^*$   
 $\int (\varphi \varphi) = (\varphi d \varphi) + f \varphi d \varphi^*$   
 $\int (\varphi \varphi) = \int (\varphi d \varphi) + f \varphi d \varphi^*$ 

### Hermitian Operators $\hat{\mathcal{A}}$

1. The eigenvalues of a Hermitian operator are real

$$\hat{A} | \psi_n \rangle = a | \psi_n \rangle$$

2. Two eigenfunctions corresponding to different eigenvalues are orthogonal

$$< \mathcal{Y}_i | \mathcal{Y}_j > = d_{ij}$$

- 3. The set of eigenfunctions of a Hermitian operator is complete  $\varphi(u) = \int \varphi(u) = \int$
- 4. Commuting Hermitian operators have a set of common eigenfunctions

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# The set of eigenfunctions of a Hermitian operator is complete



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## The set of eigenfunctions of a Hermitian operator is complete



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#### **Third Postulate**

 In any single measurement of a physical quantity that corresponds to the operator A, the only values that will be measured are the eigenvalues of that operator.

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#### Fourth Postulate

• If a series of measurements is made of the dynamical variable A on an ensemble described by  $\Psi$ , the average ("expectation") value is  $\langle A \rangle = \frac{\langle \Psi | \hat{A} | \Psi \rangle}{\langle \Psi | \Psi \rangle}$ 

i.e. the probability of obtaining an eigenvalue  $a_n$  is  $P(a_n) = |\langle \varphi_n | \Psi \rangle|^2$ 

#### **Dirac Notation**

• Eigenvalue equation:

$$\hat{A}|\psi_i\rangle = a_i|\psi_i\rangle \qquad (\Rightarrow \langle \psi_i|\psi_j\rangle = \delta_{ij})$$

• Expectation values:

$$\left\langle \psi_{i} \left| \hat{H} \psi_{i} \right\rangle = \left\langle \psi_{i} \left| \hat{H} \right| \psi_{i} \right\rangle = \int \psi_{i}^{*}(\vec{r}) \left[ -\frac{\hbar^{2}}{2m} \nabla^{2} + V(\vec{r}) \right] \psi_{i}(\vec{r}) d\vec{r} = E_{i}$$

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# Commuting Hermitian operators have a set of common eigenfunctions

#### Quantum double-slit



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#### Fifth postulate

• If the measurement of the physical quantity A gives the result  $a_n$ , the wavefunction of the system immediately after the measurement is the eigenvector  $|\varphi_n\rangle$ 

#### Position and probability



Graphs of the probability density for positions of a particle in a one-dimensional hard box according to classical mechanics removed for copyright reasons. See Mortimer, R. G. *Physical Chemistry*. 2nd ed. San Diego, CA: Elsevier, 2000, page 555, Figure 15.3.

Diagram showing the probability densities of the first 3 energy states in a 1D quantum well of width L.

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#### Quantum double-slit

Image removed due to copyright restrictions. Please see any experimental verification of the double-slit experiment, such as http://commons.wikimedia.org/wiki/Image:Doubleslitexperiment\_results\_Tanamura\_1.gif

Image of a double-slit experiment simulation removed due to copyright restrictions. Please see "Double Slit Experiment." in *Visual Quantum Mechanics*.

#### Deterministic vs. stochastic

- Classical, macroscopic objects: we have welldefined values for all dynamical variables at every instant (position, momentum, kinetic energy...)
- Quantum objects: we have well-defined probabilities of measuring a certain value for a dynamical variable, when a large number of identical, independent, identically prepared physical systems are subject to a measurement.

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#### **Top Three List**

- Albert Einstein: "Gott wurfelt nicht!" [God does not play dice!]
- Werner Heisenberg "I myself . . . only came to believe in the uncertainty relations after many pangs of conscience. . ."
- Erwin Schrödinger: "Had I known that we were not going to get rid of this damned quantum jumping, I never would have involved myself in this business!"