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3.23 Electrical, Optical, and Magnetic Properties of Materials

Fall 2007

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3.23 Fall 2007 – Lecture 5

THE HYDROGEN ECONOMY

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Last time

1. Commuting operators, Heisenberg principle
2. Measurements and collapse of the wavefunction
3. Angular momentum and spherical harmonics
4. Electron in a central potential and radial solutions

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Simultaneous eigenfunctions of L^2 , L_z

$$\hat{L}_z Y_l^m(\theta, \varphi) = m\hbar Y_l^m(\theta, \varphi)$$

$$\hat{L}^2 Y_l^m(\theta, \varphi) = \hbar^2 l(l+1) Y_l^m(\theta, \varphi)$$

$$Y_l^m(\theta, \varphi) = \Theta_l^m(\theta) \Phi_m(\varphi)$$

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An electron in a central potential

$$\hat{H} = -\frac{\hbar^2}{2\mu} \frac{1}{r^2} \frac{d}{dr} \left(r^2 \frac{d}{dr} \right) + \frac{\hat{L}^2}{2\mu r^2} + \hat{V}(r)$$

$$\psi_{nlm}(\vec{r}) = R_{nlm}(r) Y_{lm}(\vartheta, \varphi)$$

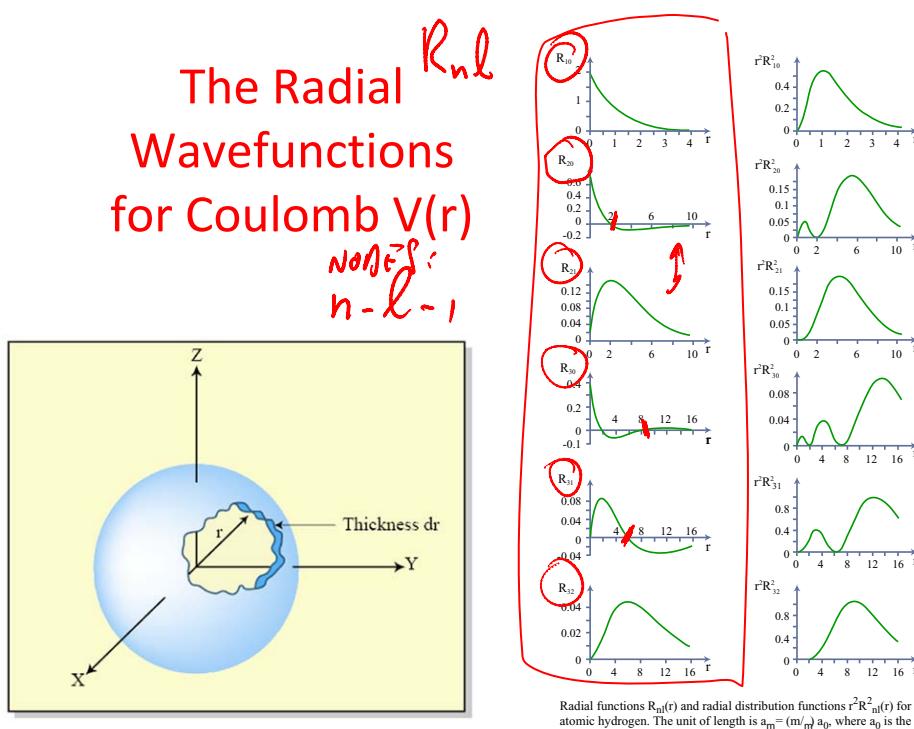
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An electron in a central potential (III)

$$u_{nl}(r) = r R_{nl}(r) \quad V_{eff}(r) = \frac{\hbar^2}{2\mu} \frac{l(l+1)}{r^2} - \frac{Ze^2}{4\pi\epsilon_0 r}$$

$$\left[-\frac{\hbar^2}{2\mu} \frac{d^2}{dr^2} + V_{eff}(r) \right] u_{nl}(r) = E_{nl} u_{nl}(r)$$

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Solutions in a Coulomb Potential

5d

4f

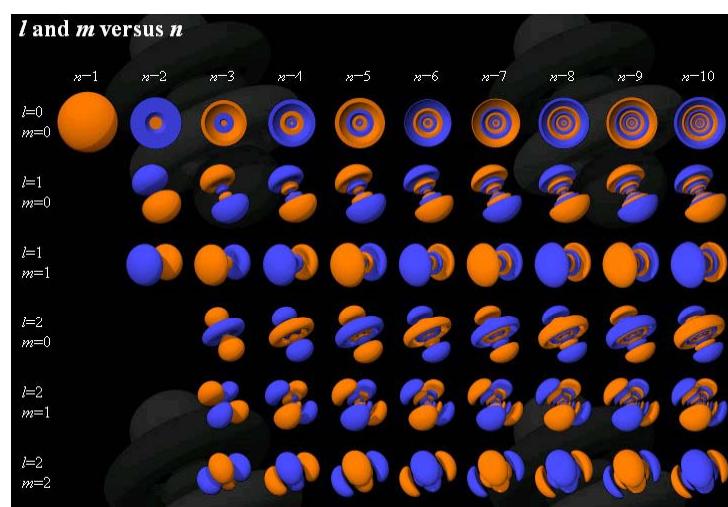
5g

Images removed; please see any visualization of the 5d, 4f, and 5g hydrogen orbitals.

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The Full Alphabet Soup

<http://www.orbitals.com/orb/orbtable.htm>



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Good Quantum Numbers

- For an operator that does not depend on t:

$$\frac{d\langle A \rangle}{dt} = \frac{d\langle \Psi | \hat{A} | \Psi \rangle}{dt} = \left\langle \frac{\partial}{\partial t} \Psi \middle| \hat{A} \middle| \Psi \right\rangle + \cancel{\left\langle \Psi \middle| \frac{\partial^2}{\partial t^2} \hat{A} \middle| \Psi \right\rangle} + \left\langle \Psi \middle| \hat{A} \middle| \frac{\partial}{\partial t} \Psi \right\rangle = \dots$$

$\hookrightarrow \frac{d}{dt} \int \Psi^* (\hat{A} \Psi) - \langle \Psi | \hat{H} - \hat{H} | \Psi \rangle = \frac{1}{i\hbar} \langle [\hat{A}, \hat{H}] \rangle$

- Then, if it commutes with the Hamiltonian, its expectation value does not change with time (it's a constant of motion – if we are in an eigenstate, that quantum number will remain constant)

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Three Quantum Numbers

- $\hat{H} \leftrightarrow$ Principal quantum number **n** (energy, accidental degeneracy)
 $E_n = -\frac{e^2}{8\pi\epsilon_0} \frac{Z^2}{a_0 n^2} = -(13.6058 \text{ eV}) \frac{Z^2}{n^2} = -(1 \text{ Ry}) \frac{Z^2}{n^2}$
ATOMIC NUMBER = 0.231 Ry
- $\hat{L}^2 \leftrightarrow$ Angular momentum quantum number **l**
 $l = 0, 1, \dots, n-1$ (a.k.a. *s, p, d... orbitals*)
- $\hat{L}_z \leftrightarrow$ Magnetic quantum number **m**
 $m = -l, -l+1, \dots, l-1, l$

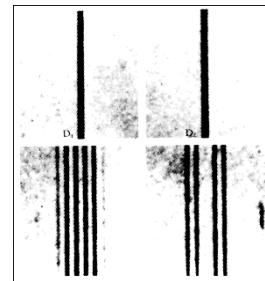
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How do you measure angular momentum ?

- Coupling to a (strong !) magnetic field \vec{B}

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Please see any experimental setup
for observing the Zeeman Effect.



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Right experiment – wrong theory (Stern-Gerlach)

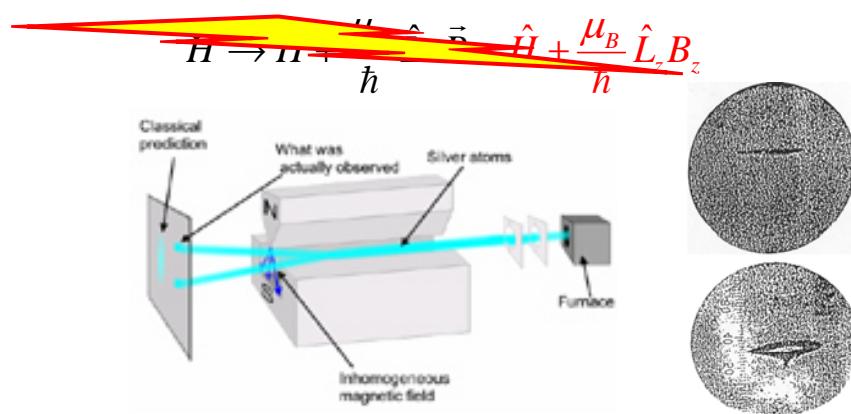


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$$\hat{H} \rightarrow \hat{H} + \frac{\mu_B}{\hbar} (\hat{L} + 2\hat{S}) \cdot \vec{B} = \hat{H} + \frac{\mu_B}{\hbar} (\hat{L}_z + 2\hat{S}_z) B_z$$

Goudsmit and Uhlenbeck

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Spin

- Dirac derived the relativistic extension of Schrödinger's equation; for a free particle he found two independent solutions for a given energy
- There is an operator (spin S) that commutes with the Hamiltonian and that can only have two eigenvalues
- In a magnetic field, the spin combines with the angular momentum, and they couple via
$$\hat{H} \rightarrow \hat{H} + \frac{\mu_B}{\hbar} (\hat{L} + 2\hat{S}) \cdot \vec{B}$$

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Spin Eigenvalues/Eigenfunctions

- Norm (s integer \rightarrow bosons, half-integer \rightarrow fermions)

$$\hat{S}^2 \Psi_{\text{spin}} = \hbar^2 s(s+1) \Psi_{\text{spin}}$$

- Z-axis projection (electron is a fermion with $s=1/2$)

$$\hat{S}_z \Psi_{\text{spin}} = \pm \frac{\hbar}{2} \Psi_{\text{spin}}$$

- Spin-orbital: product of the “space” wavefunction and the “spin” wavefunction

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Pauli Exclusion Principle

We can't have two electrons in the same quantum state →

Any two electrons in an atom cannot have the same 4 quantum numbers n, l, m, m_s

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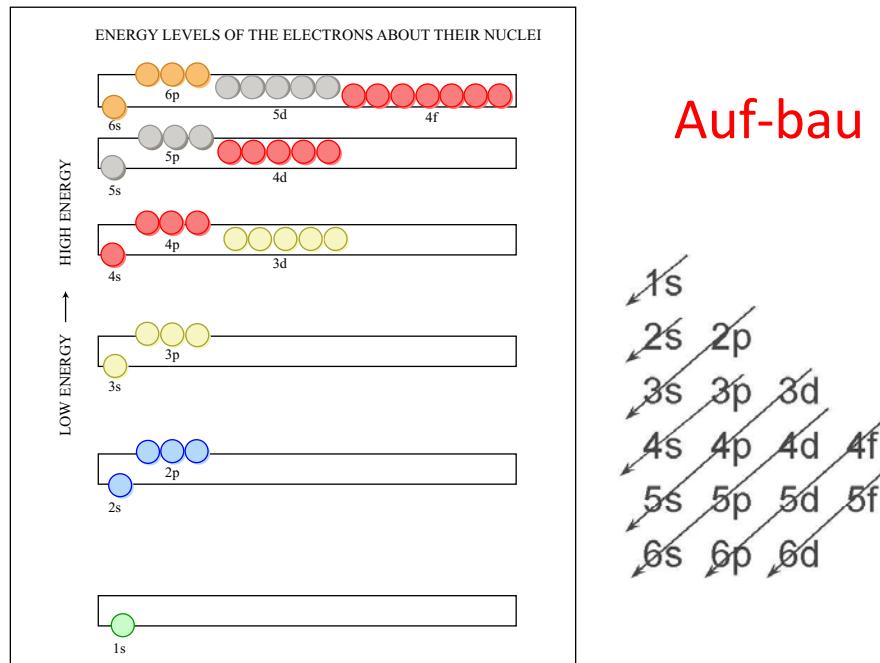


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